

The Crystal Structure of 4-Acetyl-2'-fluorobiphenyl

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The crystal structure of 4-acetyl-2'-fluorobiphenyl, $C_{14}H_{11}OF$, has been determined from three-dimensional X-ray diffraction data by a Patterson interpretation method. The crystals are monoclinic, space group $P2_1/c$, with unit-cell dimensions $a=13.687$, $b=5.971$, $c=14.766$ Å, $\beta=116^\circ 10'$. The structure was refined with use of anisotropic thermal parameters to an R value of 0.094 for the 1280 observed structure factors. The angle between the two phenyl rings was found to be 50.5° . All the intermolecular contacts correspond to normal van der Waals interactions.

Introduction

Recent studies of the liquid-crystalline properties of biphenyl derivatives (Branch, Byron, Gray, Ibbotson & Worrall, 1964) have created interest in the interplanar angles for both sterically affected and non-sterically affected biphenyl nuclei in the liquid crystalline states. Residual crystal forces must operate upon the molecules in a liquid crystal, and interplanar angles derived from studies involving solutions or the vapour phases of biphenyl derivatives are probably not relevant to the liquid-crystalline states. Thus the molecule of biphenyl is planar in the crystal (Dhar, 1932; Trotter, 1961; Hargreaves & Rizvi, 1962), but has an interplanar angle of about 45° in solution (Brande & Forbes, 1955) or in the vapour phase (Bastiansen, 1949). To help to assess more exactly the marked changes in liquid-crystalline transition temperatures which occur on introducing a suitable substituent into the 2 or 2' position of a suitable biphenyl derivative, it seemed of value to establish interplanar angles and details of molecular structure for the crystalline states of some related, but less complex, systems. Since the simple 2-halogenobiphenyls are liquids, the higher melting 4-acetyl-2'-halogenobiphenyls and 2'-halogenobiphenyl-4-carboxylic acids were selected for study (Gray, Sutherland & Young, 1965). This paper describes an X-ray crystal-

lographic study of 4-acetyl-2'-fluorobiphenyl, the molecular formula for which is shown in Fig. 1, along with the numbering system used.

The structure was also of interest as it gave an opportunity of applying the methods developed by Tollin & Cochran (1964) to a structure containing more than one planar group in the asymmetric unit.

Experimental

The crystals of 4-acetyl-2'-fluorobiphenyl, produced by slow evaporation from ethanol, were well formed bladed needles, elongated along b , the main face of the blade being the (101) plane. On examination under the polarizing microscope many of the larger crystals were found to be twinned. The space group and unit-cell dimensions were obtained from Weissenberg and precession photographs with Cu $K\alpha$ radiation. The systematic absences, $0k0$ with $k=2n+1$ and $h0l$ with $l=2n+1$, uniquely determined the space group as $P2_1/c$. The unit-cell dimensions were

$$a = 13.687 \pm 0.005, \quad b = 5.971 \pm 0.003, \quad c = 14.766 \pm 0.005 \text{ \AA};$$

$$\beta = 116^\circ 10' \pm 5'$$

$$(\lambda \text{ Cu } K\alpha_1 = 1.54050 \text{ \AA}; \lambda \text{ Cu } K\alpha_2 = 1.54434 \text{ \AA}).$$

The density measured at 17°C by the method of flotation with aqueous cadmium n-dodecatungstaborate was $1.31 \pm 0.01 \text{ g.cm}^{-3}$. The calculated value, assuming four molecules per unit cell, was 1.313 g.cm^{-3} .

In an attempt to minimize absorption effects the crystals were cut into needles of approximately square cross section. The same single crystal was used to collect the $h0l$ to $h5l$ layers by the multiple-film equi-inclination Weissenberg method with Cu $K\alpha$ radiation and the $hk0$, $0kl$, hkh , $h,k,2h$, $hk\bar{h}$ and $h,k,2\bar{h}$ spectra on a precession camera using timed exposures with Cu $K\alpha$ and Mo $K\alpha$ radiation. The intensity of every

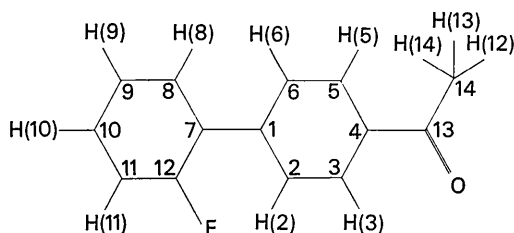


Fig. 1. The molecular formula of 4-acetyl-2'-fluorobiphenyl and the numbering adopted throughout the paper.

spot on all films was estimated visually by comparison with a calibrated scale and corrected for Lorentz and polarization factors. Corrections for spot shape were applied to the upper level Weissenberg data: the Phillips (1954) correction factor to the extended reflexions and an empirical correction to the contracted reflexions obtained by a comparison of the intensity of the contracted reflexion with that of the corrected extended reflexion, for those reflexions common to both sides of the film. The intensities from the precession films, scaled to the $h0l$ data, were used to increase the number of such comparisons at low values of $\sin \theta$. After the data were brought on to the same arbitrary scale, using the common reflexions and timed exposures, and before averaging, the agreement between different observations of the same structure factor was of the order of 10%.

Structure determination

The function $I(\theta, \varphi)$ (Tollin & Cochran, 1964) was computed from the 50 largest sharpened Patterson coefficients, $|F_s(\mathbf{h})|^2$, and $R=3.4 \text{ \AA}$ and plotted on a Sanson-Flamsteed projection (Fig. 2). The projection contained only one large peak which had a little fine structure suggesting that there might be two very close peaks which overlap. The $I(\theta, \varphi)$ map was recomputed in the region of this peak with the use of all the $|F_s|^2$ data. This showed a single peak from which the fine structure had disappeared. The angular coordinates of the peak were $\theta=27.5^\circ$ and $\varphi=100^\circ$ where θ and φ are spherical polar angles referred to the orthogonal triad $\mathbf{a}^*, \mathbf{b}^*, \mathbf{c}$.

This single peak can be interpreted as indicating that either the whole molecule is essentially planar, or that the angle between the two planar parts of the molecule is such that the peak due to one part overlaps the peak

due to the second part in a symmetry related molecule. This would occur if the normals to the two planar parts were related by the equations $\theta_1=\theta_2$ and $\varphi_1=-\varphi_2$. These two possibilities were denoted as the parallel and anti-parallel configurations.

The section through the sharpened Patterson function, computed in the orientation given by these θ and φ values, is shown in Fig. 3. The line X is the intersection of the molecular plane with the (010) plane.

Theoretical vector sets, constructed for both the parallel and anti-parallel configurations of the biphenyl moiety and the atom C(13), were superimposed on the Patterson section and rotated until the best fit was obtained. It was impossible to distinguish between the two configurations; both gave the best fit when oriented so that the long axis of the molecule lay along the line L in Fig. 3. The orientation of the acetyl group and the position of the fluorine atom could not be identified at this stage.

To distinguish between the parallel and antiparallel configurations and determine the position of the molecule in the unit cell, the method of Q -functions (Tollin, 1966) was used. The function $Q(Y_0)$ which determines the position of the molecule relative to the c -glide was

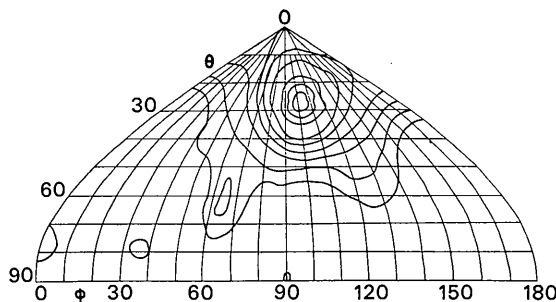


Fig. 2. Plot of the $I(\theta, \varphi)$ function.

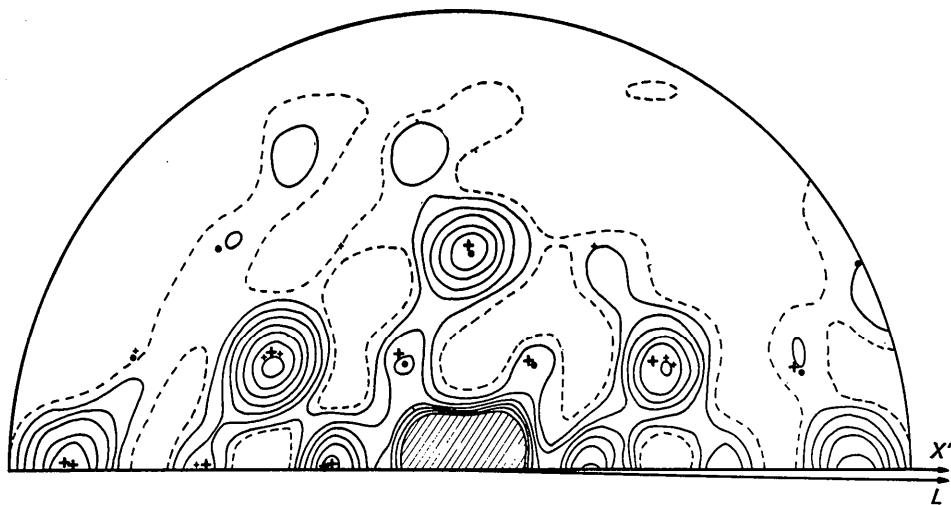


Fig. 3. Section through the sharpened Patterson function, with contours at equal arbitrary intervals. Superimposed is the theoretical vector set, the relative weights of the vectors are shown by black dots for weight 1, small + for weight 2 and large + for weight 4. The radius of the section is 5 \AA .

calculated for both configurations (Fig. 4). In Fig. 4(a), corresponding to the parallel configuration, all three peaks are approximately the same height, whereas in Fig. 4(b) corresponding to the antiparallel configuration there is one peak which is larger than the other two. The function $Q(X_0, Z_0)$ was calculated for both configurations to determine the position of the molecule relative to the Z_1 axis. The map for the antiparallel configuration is shown in Fig. 5. The centre of symmetry in this map, and the centre in the $Q(Y_0)$ function for the parallel configuration, arise from the symmetrical shape of the part of the molecule used. Within the asymmetric unit of the map in Fig. 5 there is one peak which is very large compared with the remainder whereas the function $Q(X_0, Z_0)$ for the parallel configuration showed no outstandingly large peak. It was concluded that the antiparallel configuration was the correct one, and that the position of the molecule in the unit cell was defined by the largest peaks in Fig. 4(b) and Fig. 5.

There are two possible positions for the acetyl group, assuming it to be coplanar with the phenyl ring to which it is attached. On the basis of structure factor calculations on the $(h0l)$ zone, one of the possibilities was eliminated. A similar technique was employed to position the fluorine atom. The structure was then refined in projection to $R(h0l)=0.14$, where $R = \frac{\sum ||F_o| - |F_c||}{\sum |F_o|}$ initially by Fourier refinement and then by a least-squares process using an overall isotropic temperature factor. The least-squares refinement program used was that due to J. C. Schoone (Rogers, 1961) for the Statec ZEBRA computer.

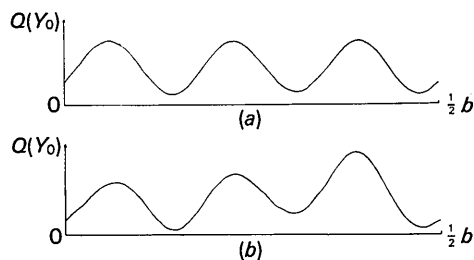


Fig. 4. The function $Q(Y_0)$, (a) for the parallel configuration, (b) for the anti-parallel configuration.

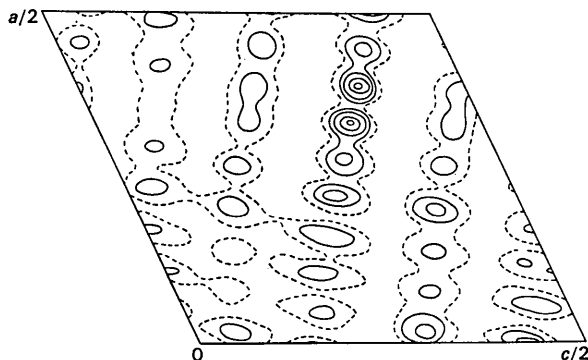


Fig. 5. The function $Q(X_0, Z_0)$ for anti-parallel configuration.

In order to confirm the proposed structure, a refinement based on the refined x and z coordinates and the y coordinates from $Q(Y_0)$ was carried out on the 460 largest general structure factors. The structure was quickly refined to $R=0.15$.

Refinement

Refinement of the structure was continued with all the observed data, by least-squares using an overall isotropic temperature factor and proceeded from $R=0.35$ to 0.19 in 5 cycles. A three-dimensional electron density map was computed with the program of Daly, Stephens & Wheatley (1963) for the Elliot 803B computer, to check that the peak heights and the background were reasonable. Introducing individual isotropic temperature factors a further four cycles reduced R to 0.17 . Coordinates for eight of the eleven hydrogen atoms were computed on the assumption of their planarity with the phenyl ring to which they were attached and a C-H bond length of 1.05 \AA . After two further cycles R fell to 0.15 . Further least-squares refinement was carried out with the SFLS program of G. A. Mair for the IBM 1620 II computer. The program uses the block-diagonal approximation and the coefficients were weighted by

$$\omega = 1 / \{1 + (KF_o - b)^2 / a^2\},$$

where a and b are constants selected to discriminate against weak reflexions and against very strong reflexions. Anisotropic temperature factors of the form

$$\exp [-(B_{11}h^2 + B_{22}k^2 + B_{33}l^2 + B_{12}hk + B_{23}kl + B_{13}hl)]$$

were incorporated into the refinement process which proceeded to $R=0.11$. A three-dimensional difference Fourier synthesis showed regions of density up to 0.4 e. \AA^{-3} around the methyl carbon atom. These regions were in the form of long elongated distributions curving round the carbon atom, suggesting that the methyl group may undergo hindered rotation about the C-C bond. Hydrogen atom positions were allocated to three of the highest regions, lying between 0.9 and 1.2 \AA from the carbon atom and chosen so that the C-C-H bond angles were of the order of 110° . Further refinement proceeded to $R=0.094$ for the observed data. In the latter stages of refinement the reflexions 002 and 204 were excluded as it was suspected that these suffered from extinction. The temperature parameters of the methyl hydrogen atoms increased to about 8 \AA^2 , much larger than the B values for the hydrogen atoms connected to the phenyl rings and consistent with the smearing effect noted in the difference synthesis. Little significance can be attributed to the positions of the methyl hydrogen atoms. The final structure factors are given in Table 1. Table 2 gives the final heavy-atom coordinates and anisotropic thermal parameters. The hydrogen atom parameters appear in Table 3.

Table 1. The observed and calculated structure factors

h	k	l	observed	calculated
0	0	0	1000	1000
0	0	1	300	300
0	0	2	376	376
0	0	3	166	166
0	0	4	289	289
0	0	5	163	163
0	0	6	329	329
0	0	7	199	199
0	0	8	306	306
0	0	9	187	187
0	0	10	324	324
0	0	11	199	199
0	0	12	306	306
0	0	13	187	187
0	0	14	324	324
0	0	15	199	199
0	0	16	306	306
0	0	17	187	187
0	0	18	324	324
0	0	19	199	199
0	0	20	306	306
0	0	21	187	187
0	0	22	324	324
0	0	23	199	199
0	0	24	306	306
0	0	25	187	187
0	0	26	324	324
0	0	27	199	199
0	0	28	306	306
0	0	29	187	187
0	0	30	324	324
0	0	31	199	199
0	0	32	306	306
0	0	33	187	187
0	0	34	324	324
0	0	35	199	199
0	0	36	306	306
0	0	37	187	187
0	0	38	324	324
0	0	39	199	199
0	0	40	306	306
0	0	41	187	187
0	0	42	324	324
0	0	43	199	199
0	0	44	306	306
0	0	45	187	187
0	0	46	324	324
0	0	47	199	199
0	0	48	306	306
0	0	49	187	187
0	0	50	324	324
0	0	51	199	199
0	0	52	306	306
0	0	53	187	187
0	0	54	324	324
0	0	55	199	199
0	0	56	306	306
0	0	57	187	187
0	0	58	324	324
0	0	59	199	199
0	0	60	306	306
0	0	61	187	187
0	0	62	324	324
0	0	63	199	199
0	0	64	306	306
0	0	65	187	187
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0	0	67	199	199
0	0	68	306	306
0	0	69	187	187
0	0	70	324	324
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0	0	78	324	324
0	0	79	199	199
0	0	80	306	306
0	0	81	187	187
0	0	82	324	324
0	0	83	199	199
0	0	84	306	306
0	0	85	187	187
0	0	86	324	324
0	0	87	199	199
0	0	88	306	306
0	0	89	187	187
0	0	90	324	324
0	0	91	199	199
0	0	92	306	306
0	0	93	187	187
0	0	94	324	324
0	0	95	199	199
0	0	96	306	306
0	0	97	187	187
0	0	98	324	324
0	0	99	199	199
0	0	100	306	306
0	0	101	187	187
0	0	102	324	324
0	0	103	199	199
0	0	104	306	306
0	0	105	187	187
0	0	106	324	324
0	0	107	199	199
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0	0	111	199	199
0	0	112	306	306
0	0	113	187	187
0	0	114	324	324
0	0	115	199	199
0	0	116	306	306
0	0	117	187	187
0	0	118	324	324
0	0	119	199	199
0	0	120	306	306
0	0	121	187	187
0	0	122	324	324
0	0	123	199	199
0	0	124	306	306
0	0	125	187	187
0	0	126	324	324
0	0	127	199	199
0	0	128	306	306
0	0	129	187	187
0	0	130	324	324
0	0	131	199	199
0	0	132	306	306
0	0	133	187	187
0	0	134	324	324
0	0	135	199	199
0	0	136	306	306
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0	0	182	324	324
0	0	183	199	199
0	0	184	306	306
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0	0	186	324	324
0	0	187	199	199
0	0	188	306	306
0	0	189	187	187
0	0	190	324	324
0	0	191	199	199
0	0	192	306	306
0	0	193	187	187
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0	0	196	306	306
0	0	197	187	187
0	0	198	324	324
0	0	199	199	199
0	0	200	306	306
0	0	201	187	187
0	0	202	324	324
0	0	203	199	199
0	0	204	306	306
0	0	205	187	187
0	0	206	324	324
0	0	207	199	199
0	0	208	306	306
0	0	209	187	187
0	0	210	324	324
0	0	211	199	199
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0	0	256	306	306
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0	0	258	324	324
0	0	259	199	199
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0	0	261	187	187
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0	0	290	324	324
0	0	291	199	199
0	0	292	306	306
0	0	293	187	187
0	0	294	324	324
0	0	295	199	199
0	0	296	306	306
0	0	297	187	187
0	0	298	324	324

Table 2. Final fractional coordinates and anisotropic temperature factors for the non-hydrogen atoms

The anisotropic temperature factors have been multiplied by 10^5 .

	<i>x</i>	<i>y</i>	<i>z</i>	<i>B</i> ₁₁	<i>B</i> ₂₂	<i>B</i> ₃₃	<i>B</i> ₂₃	<i>B</i> ₃₁	<i>B</i> ₁₂
F	0.3757	0.8924	0.4910	782	1815	787	1043	603	-461
O	-0.1636	0.7933	0.3252	791	4514	1251	464	658	1501
C(1)	0.2287	0.5682	0.4937	536	1234	465	194	452	-26
C(2)	0.1686	0.3891	0.4340	562	1244	643	-83	327	313
C(3)	0.0568	0.4046	0.3771	679	2041	511	-196	215	-235
C(4)	0.0018	0.5981	0.3809	627	2091	461	463	466	381
C(5)	0.0609	0.7719	0.4398	698	1875	727	205	755	546
C(6)	0.1704	0.7592	0.4959	727	1669	531	-641	528	-519
C(7)	0.3478	0.5501	0.5550	580	1077	438	-59	476	-220
C(8)	0.3953	0.3651	0.6171	474	1780	489	201	257	-328
C(9)	0.5069	0.3463	0.6701	740	2084	549	342	523	230
C(10)	0.5745	0.5131	0.6648	503	2994	610	-8	466	-136
C(11)	0.5288	0.6970	0.6042	716	2261	612	-383	628	-796
C(12)	0.4190	0.7108	0.5525	631	1537	449	152	403	73
C(13)	-0.1197	0.6197	0.3216	687	3811	504	478	424	173
C(14)	-0.1857	0.4274	0.2603	597	5594	793	322	214	-121

Discussion

The arrangement of the molecules in the unit cell viewed along the *b* axis is shown in Fig. 6. The intramolecular bond lengths and bond angles along with their standard deviations are given in Table 4.

Table 3. The hydrogen atom fractional coordinates

	<i>x</i>	<i>y</i>	<i>z</i>
H(2)	0.209	0.253	0.429
H(3)	0.011	0.277	0.335
H(5)	0.026	0.894	0.439
H(6)	0.212	0.878	0.528
H(8)	0.342	0.229	0.618
H(9)	0.530	0.246	0.709
H(10)	0.642	0.508	0.692
H(11)	0.573	0.826	0.595
H(12)	-0.269	0.446	0.242
H(13)	-0.180	0.417	0.195
H(14)	-0.171	0.282	0.291

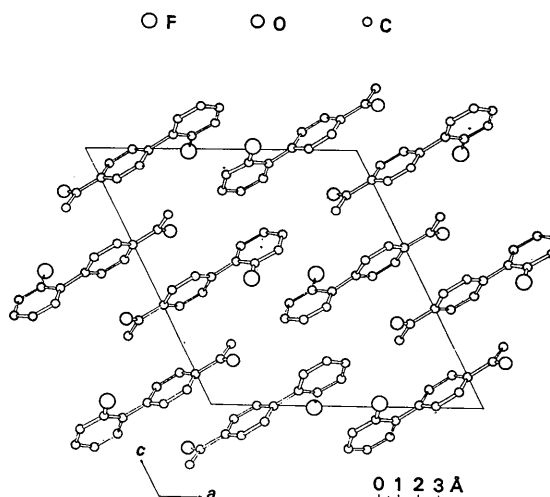
Fig. 6. The arrangement of the molecules viewed along the *b* axis.

Table 4. The bond lengths and bond angles with their estimated standard deviations

	Length	e.s.d.		Length	e.s.d.
F—C(12)	1.369 Å	0.008 Å	C(5)—C(6)	1.358 Å	0.012 Å
O—C(13)	1.211	0.009	C(7)—C(8)	1.399	0.009
C(1)—C(2)	1.399	0.009	C(7)—C(12)	1.380	0.008
C(1)—C(6)	1.400	0.008	C(8)—C(9)	1.380	0.011
C(1)—C(7)	1.479	0.010	C(9)—C(10)	1.385	0.009
C(2)—C(3)	1.386	0.012	C(10)—C(11)	1.380	0.010
C(3)—C(4)	1.394	0.009	C(11)—C(12)	1.355	0.011
C(4)—C(5)	1.366	0.010	C(13)—C(14)	1.495	0.013
C(4)—C(13)	1.505	0.012			
	Angle	e.s.d.		Angle	e.s.d.
C(2)—C(1)—C(6)	117.0°	0.7°	C(8)—C(7)—C(12)	116.0°	0.7°
C(2)—C(1)—C(7)	120.8	0.5	C(7)—C(8)—C(9)	120.9	0.5
C(6)—C(1)—C(7)	122.1	0.6	C(8)—C(9)—C(10)	120.7	0.6
C(1)—C(2)—C(3)	121.1	0.5	C(9)—C(10)—C(11)	119.1	0.7
C(2)—C(3)—C(4)	120.2	0.7	C(10)—C(11)—C(12)	119.1	0.6
C(3)—C(4)—C(5)	118.5	0.7	F—C(12)—C(7)	117.7	0.7
C(3)—C(4)—C(13)	121.7	0.7	F—C(12)—C(11)	118.0	0.5
C(5)—C(4)—C(13)	119.8	0.6	C(7)—C(12)—C(11)	124.3	0.6
C(4)—C(5)—C(6)	122.0	0.6	O—C(13)—C(4)	119.3	0.7
C(1)—C(6)—C(5)	121.2	0.6	O—C(13)—C(14)	120.4	0.8
C(1)—C(7)—C(8)	121.8	0.5	C(4)—C(13)—C(14)	120.3	0.6
C(1)—C(7)—C(12)	122.2	0.6			

The equations for the mean plane of the ring systems, C(1)–C(6): C(7)–C(12): C(7)–C(12) and F and the acetyl radical calculated by the method of Schomaker, Waser, Marsh & Bergman (1959), are

$$\begin{aligned} -0.5182X - 0.3746Y + 0.7689Z &= 3.8148 \\ -0.4239X + 0.4724Y + 0.7728Z &= 6.7559 \\ -0.4249X + 0.4784Y + 0.7685Z &= 6.7350 \\ -0.5049X - 0.3404Y + 0.7934Z &= 4.0079 \end{aligned}$$

where X, Y, Z are the atomic coordinates in Å referred to the orthogonal axial triad a, b, c^* . The deviations of the atoms from the mean planes of the phenyl rings are given in Table 5. The angle between the acetyl radical and the phenyl ring C(1)–C(6) is 2.5° and the angle between the two phenyl rings is 50.5° . The equivalent angle calculated from the $I(\theta, \varphi)$ map was 55° . The normal to the molecular plane as given by the $I(\theta, \varphi)$ function and the normals to the phenyl rings are thus in agreement to within 2.4° .

Table 5. *The deviations of the atoms from the mean planes of the phenyl rings*

C(1)	-0.009 Å	C(7)	-0.005 Å
C(2)	+0.008	C(8)	+0.002
C(3)	-0.005	C(9)	-0.002
C(4)	+0.004	C(10)	+0.003
C(5)	-0.005	C(11)	-0.006
C(6)	+0.008	C(12)	+0.006

The distance of 1.369 \AA for the carbon–fluorine bond, obtained in this analysis, is considerably larger than the value published in *International Tables for X-ray Crystallography* (1962) of $1.30 \pm 0.01 \text{ \AA}$. In *o*-fluorobenzoic acid (Krause & Dunken, 1966) the C–F distance was determined as $1.368 \pm 0.019 \text{ \AA}$ in good agreement with the value found in the present analysis. The bond length distribution found in the fluoro-substituted phenyl ring is remarkably similar to that in *o*-fluorobenzoic acid, as can be seen in Table 6, which also shows the averaged equivalent bond lengths in biphenyl.

Table 6. *Comparison of equivalent bond distances in the fluoro-substituted phenyl ring in 4-acetyl-2'-fluorobiphenyl (YTS) and o-fluorobenzoic acid (KD) with those in the phenyl ring of biphenyl (T, HR)**

Bond	(YTS)	(KD)	(T)	(HR)
C(7)–C(12)	1.380	1.380	1.372	1.383
C(12)–C(11)	1.355	1.365	1.416	1.421
C(11)–C(10)	1.380	1.395	1.376	1.361
C(10)–C(9)	1.385	1.395	1.376	1.361
C(9)–C(8)	1.380	1.375	1.416	1.421
C(8)–C(7)	1.399	1.395	1.372	1.383
Mean	1.382	1.384		

* YTS this investigation
 KD Krause & Dunken (1966)
 T Trotter (1961)
 HR Hargreaves & Rizvi (1962)

The C(1)–C(7) bond length of $1.479 \pm 0.010 \text{ \AA}$ is significantly shorter than the quoted value of 1.54 \AA . However in biphenyl this measured $1.506 \pm 0.017 \text{ \AA}$, which is possibly significantly larger than in this determination. This is perhaps due to the fact that biphenyl

is a planar molecule and the steric repulsion between H(2) and H(12) and between H(6) and H(8) causes a strain between C(1) and C(7) which lengthens the bond. In the present structure determination, the molecule is non-planar and the steric repulsion between H(6) and H(8) no longer exists because of the twist about C(1)–C(7). The F–H(2) distance of 2.53 \AA is just smaller than the value of 2.55 \AA obtained by considering the van der Waals radii. The value obtained for C(1)–C(7) is in good agreement with the theoretical calculations of Dewar & Schmeising (1959) and Cruickshank & Sparks (1960), who suggest that the (sp^2) single-bond length between trigonally linked carbon atoms should be taken as 1.477 \AA . This would indicate that there was no conjugation taking place between the rings in this structure.

All the intermolecular distances correspond to normal van der Waals interactions. All contacts less than 3.6 \AA are given in Table 7.

Table 7. *Intermolecular distances between atoms less than 3.6 \AA apart (excluding hydrogen atoms)*

i	j	Equipoint indication	d_{ij}
F	C(8)	$x, 1+y, z$	3.326 \AA
C(10)	O	$1+x, 1.5-y, 0.5+z$	3.508
C(11)	F	$1-x, 2-y, 1-z$	3.364
C(12)	C(10)	$1-x, 1-y, 1-z$	3.513
C(14)	F	$-x, -0.5+y, 0.5-z$	3.473

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References

- BASTIANSEN, O. (1949). *Acta Chem. Scand.* **3**, 408.
 BRANCH, S. G., BYRON, D. J., GRAY, G. W., IBBOTSON, A. & WORRALL, B. M. (1964). *J. Chem. Soc.* p. 3279.
 BRANDE, E. A. & FORBES, W. F. (1955). *J. Chem. Soc.* p. 3777.
 CRUICKSHANK, D. W. J. & SPARKS, R. A. (1960). *Proc. Roy. Soc. A*, **258**, 270.
 DALY, J. J., STEPHENS, F. S. & WHEATLEY, P. J. (1963). M.R.S.A. Final Report no. 52.
 DEWAR, M. J. S. & SCHMEISING, H. N. (1959). *Tetrahedron*, **5**, 166.
 DHAR, J. (1932). *Indian J. Phys.* **7**, 43.
 GRAY, G. W., SUTHERLAND, H. H. & YOUNG, D. W. (1965). *J. Chem. Soc.* p. 4208.
 HARGREAVES, A. & RIZVI, S. H. (1962). *Acta Cryst.* **15**, 365. *International Tables for X-ray Crystallography* (1962). Vol. III. Birmingham: Kynoch Press.
 KRAUSSE, J. & DUNKEN, H. (1966). *Acta Cryst.* **20**, 67.
 PHILLIPS, D. C. (1954). *Acta Cryst.* **7**, 746.
 ROGERS, D. (1961). In *Computing Methods and the Phase Problem in X-ray Crystal Analysis*. Oxford: Pergamon Press.
 SCHOMAKER, V., WASER, J., MARSH, R. E. & BERGMAN, G. (1959). *Acta Cryst.* **12**, 600.
 TOLLIN, P. & COCHRAN, W. (1964). *Acta Cryst.* **17**, 1322.
 TOLLIN, P. (1966). *Acta Cryst.* **21**, 613.
 TROTTER, J. (1961). *Acta Cryst.* **14**, 1135.